

# PHYS 461: Final Exam

1 From the relations

$$\frac{\chi_T}{\chi_S} = \frac{C_P}{C_V}, \quad \text{and} \quad C_P - C_V = \frac{TV\alpha^2}{\chi_T},$$

where  $C_V$  is the heat capacity at constant volume, it follows that

$$\begin{aligned} \chi_T - \chi_S &= \chi_T \left( 1 - \frac{\chi_S}{\chi_T} \right) \\ &= \chi_T \left( 1 - \frac{C_V}{C_P} \right) \\ &= \chi_T \left( \frac{C_P - C_V}{C_P} \right) \\ &= \frac{TV\alpha^2}{C_P}. \end{aligned}$$

2 (i) Since  $n$  fish are marked among  $N$ ,

- $\binom{N}{n}$  is the number of different ways to choose the  $n$  fish to be marked among  $N$ .
- $\binom{n}{k}$  is the number of different ways to choose  $k$  marked fish among  $n$ .
- $\binom{N-n}{n-k}$  is the number of different ways to choose the  $n-k$  fish among the  $N-n$  non marked ones.

Hence,

$$p(N, n, k) = \frac{\binom{n}{k} \binom{N-n}{n-k}}{\binom{N}{n}}.$$

$2n - k$  being the number of different fish caught, this result makes sense only if  $N > 2n - k$ .

Consider

$$\begin{aligned}\frac{p(N, n, k)}{p(N-1, n, k)} &= \frac{\binom{n}{k} \binom{N-n}{n-k} \binom{N-1}{n}}{\binom{n}{k} \binom{N-1-n}{n-k} \binom{N}{n}} \\ &= \frac{(N-n)^2}{N(N-2n+k)}.\end{aligned}$$

From

$$\begin{aligned}\frac{(N-n)^2}{N(N-2n+k)} > 1 &\Rightarrow n^2 > kN \\ \frac{(N-n)^2}{N(N-2n+k)} < 1 &\Rightarrow n^2 < kN\end{aligned}$$

Hence, as a function of  $N$ ,  $p(N, n, k)$  increases for  $N < n^2/k$ , and decreases for  $N > n^2/k$ . The probability  $p(N, n, k)$  is, therefore, maximum when  $N$  is the largest integer less than or equal to  $n^2/k$ .

That is, the *maximum likelihood estimate* of  $N$  is given by

$$N_{\max} = \left\lfloor \frac{n^2}{k} \right\rfloor,$$

where  $\lfloor x \rfloor$  is largest integer less than or equal to  $x$ .

If  $n = 1000$  and  $k = 100$ , the maximum likelihood estimate of the number of fish is  $N_{\max} = 10000$ .

**3** (i) The partition function  $Z(N, \beta, B)$  of a  $N$  noninteracting spins  $\frac{1}{2}$  in an external magnetic field  $B$  directed along the  $z$ -axis is given by the sum over all the values of the  $z$ -components of all the spins  $S^i$ . Since

$$\exp\left(\beta g \mu B \sum_{i=1}^N S_z^i\right) = \prod_{i=1}^N \exp(\beta g \mu B S_z^i),$$

the partition function is given by

$$\begin{aligned} Z(\beta, B, N) &= \left(\exp\left(\frac{1}{2}\beta g \mu B\right) + \exp\left(-\frac{1}{2}\beta g \mu B\right)\right)^N \\ &= \left(2 \cosh\left(\frac{1}{2}\beta g \mu B\right)\right)^N. \end{aligned}$$

(ii) The Gibbs potential is given by

$$\begin{aligned} G(\beta, B, N) &= -\frac{1}{\beta} \log Z(\beta, B, N) \\ &= -\frac{N}{\beta} \log \left(2 \cosh\left(\frac{1}{2}\beta g \mu B\right)\right). \end{aligned}$$

(iii) The entropy is given by

$$S(\beta, B, N) = -\frac{\partial G}{\partial \beta} \frac{d\beta}{dT} = k\beta^2 \frac{\partial G}{\partial \beta}.$$

Hence

$$S(\beta, B, N) = kN \log \left( 2 \cosh\left(\frac{1}{2}\beta g\mu B\right) \right) \\ - \frac{1}{2} kN \beta g\mu B \tanh\left(\frac{1}{2}\beta g\mu B\right).$$

(iv) The specific heat per spin at constant field is defined as

$$c_B(\beta, B) = \frac{T}{N} \frac{\partial S}{\partial T} = -\frac{\beta}{N} \frac{\partial S}{\partial \beta},$$

that is,

$$c_B(\beta, B) = \frac{k\beta^2 g^2 \mu^2 B^2}{4 \cosh^2\left(\frac{1}{2}\beta g\mu B\right)}.$$

This quantity is obviously positive.

(v) When, for a given value of the applied field  $B$ ,  $\beta$  tends to infinity, the hyperbolic cosine grows much faster than  $\beta^2$ , thus  $c_B(\beta, B)$  tends to zero. Since  $\beta \rightarrow \infty$  means  $T \rightarrow 0$ , this result agrees with the third principle of thermodynamics.